

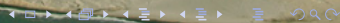
Non-homogeneous incompressible Bingham flows with variable yield stress and application to volcanology.

Jordane MATHÉ
with Laurent CHUPIN (maths) and
Karim KELFOUN (LMV)

Laboratoire de Maths & Laboratoire Magmas et Volcans

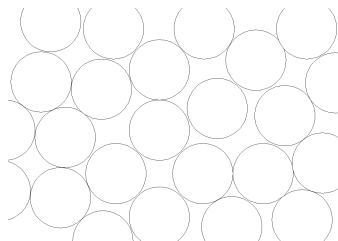
june 2015





Laboratory experiment

Zoom into the granular column \rightarrow

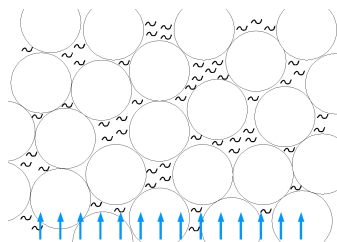


O. ROCHE, LMV.

Laboratory experiment

Zoom into the granular column →

Fluidisation :
injection of **gas**
through
a pore plate.



O. ROCHE, LMV.

1 Model

2 Numerical simulation

3 Perspectives

1 Model



Two phases for one fluid

Consider **one mixed fluid** with **variable density**.

Two phases for one fluid

Consider **one mixed fluid** with **variable density**. It satisfies the non-homogeneous incompressible Navier-Stokes equations:

$$\begin{cases} \operatorname{div}(\mathbf{v}) = 0 \\ \partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = 0 \\ \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) + \nabla p = \rho \mathbf{g} + \operatorname{div}(\mathbf{S}) \end{cases}$$

Two phases for one fluid

Consider **one mixed fluid** with **variable density**. It satisfies the non-homogeneous incompressible Navier-Stokes equations:

$$\begin{cases} \operatorname{div}(\mathbf{v}) = 0 \\ \partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = 0 \\ \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) + \nabla p = \rho \mathbf{g} + \operatorname{div}(\mathbf{S}) \end{cases}$$

Unknown:

- \mathbf{v} : velocity,
- p : total pressure,
- ρ : density.

Two phases for one fluid

Consider **one mixed fluid** with **variable density**. It satisfies the non-homogeneous incompressible Navier-Stokes equations:

$$\begin{cases} \operatorname{div}(\boldsymbol{v}) = 0 \\ \partial_t \rho + \operatorname{div}(\rho \boldsymbol{v}) = 0 \\ \partial_t(\rho \boldsymbol{v}) + \operatorname{div}(\rho \boldsymbol{v} \otimes \boldsymbol{v}) + \nabla p = \rho \boldsymbol{g} + \operatorname{div}(\mathbf{S}) \end{cases}$$

Constant

- \boldsymbol{g} : gravity.

Two phases for one fluid

Consider **one mixed fluid** with **variable density**. It satisfies the non-homogeneous incompressible Navier-Stokes equations:

$$\begin{cases} \operatorname{div}(\boldsymbol{v}) = 0 \\ \partial_t \rho + \operatorname{div}(\rho \boldsymbol{v}) = 0 \\ \partial_t(\rho \boldsymbol{v}) + \operatorname{div}(\rho \boldsymbol{v} \otimes \boldsymbol{v}) + \nabla p = \rho \boldsymbol{g} + \operatorname{div}(\mathbf{S}) \end{cases}$$

Constant

- \boldsymbol{g} : gravity.

Rheology

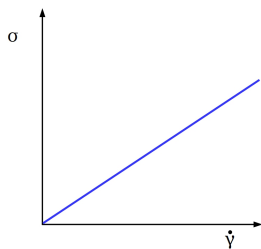
Let precise \mathbf{S} .

Rheology

In our case

$$\mathbf{S} = \mu D\mathbf{v} +$$

where $D\mathbf{v} = \dot{\gamma}$ is the strain rate tensor,
 μ is the effective viscosity



Rheology

In our case

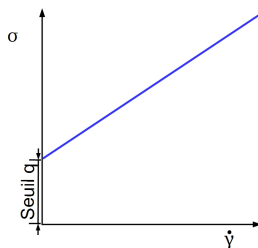
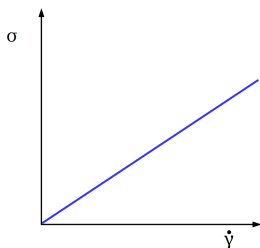
$$\mathbf{S} = \mu D\mathbf{v} + \Sigma,$$

where

$D\mathbf{v} = \dot{\gamma}$ is the strain rate tensor,
 μ is the effective viscosity

$$\Sigma = q \frac{D\mathbf{v}}{|D\mathbf{v}|}$$

\Rightarrow Bingham fluid with yield stress q



Rheology

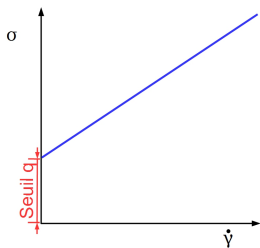
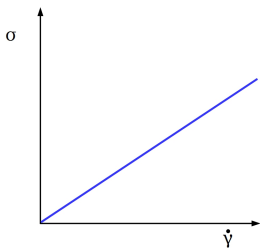
In our case

$$\mathbf{S} = \mu D\mathbf{v} + \Sigma,$$

where $D\mathbf{v} = \dot{\gamma}$ is the strain rate tensor,
 μ is the effective viscosity

$$\Sigma = q \frac{D\mathbf{v}}{|D\mathbf{v}|}$$

⇒ Bingham fluid with yield stress q



Idea

Let vary the yield stress q as a function of the **interstitial gas pressure**.

Variation of the yield stress

Definition of the yield stress

$$q = \begin{cases} \text{atmospheric pressure:} & \text{Coulomb friction} \\ \text{high pressure:} & \text{fluid} \end{cases}$$

Variation of the yield stress

Definition of the yield stress

$$q = \begin{cases} \text{atmospheric pressure:} & \text{Coulomb friction} \\ \text{high pressure:} & \text{fluid} \end{cases}$$
$$q = \begin{cases} \tan(\delta)(\rho gh - \text{gas pressure}) & \text{if low gas pressure} \\ 0 & \text{if high gas pressure} \end{cases}$$

where δ is the internal friction angle.

Variation of the yield stress

Definition of the yield stress

$$\begin{aligned} q &= \begin{cases} \text{atmospheric pressure:} & \text{Coulomb friction} \\ \text{high pressure:} & \text{fluid} \end{cases} \\ q &= \begin{cases} \tan(\delta)(\rho gh - \text{gas pressure}) & \text{if low gas pressure} \\ 0 & \text{if high gas pressure} \end{cases} \\ q &= \tan(\delta)(\rho gh - \text{gas pressure})^+ \end{aligned}$$

where δ is the internal friction angle.

Equations of the model

Noting p_f the interstitial gas pressure, we obtain:


$$\left\{ \begin{array}{l} \operatorname{div}(v) = 0 \\ \partial_t \rho + \operatorname{div}(\rho v) = 0 \\ \partial_t(\rho v) + \operatorname{div}(\rho v \otimes v) - \mu \Delta v + \nabla p = \tan(\delta) \operatorname{div} \left(\underbrace{(\rho g h - p_f)^+}_{\text{yield } q} \frac{Dv}{|Dv|} \right) + \rho \mathbf{g} \end{array} \right.$$

Equations of the model

Noting p_f the interstitial gas pressure, we obtain:

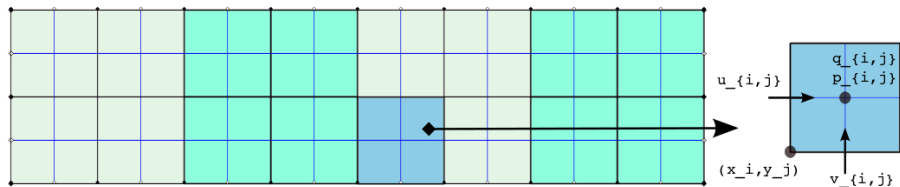
$$\left\{ \begin{array}{l} \operatorname{div}(v) = 0 \\ \partial_t \rho + \operatorname{div}(\rho v) = 0 \\ \partial_t(\rho v) + \operatorname{div}(\rho v \otimes v) - \mu \Delta v + \nabla p = \tan(\delta) \operatorname{div} \left(\underbrace{(\rho g h - p_f)^+}_{\text{yield } q} \frac{Dv}{|Dv|} \right) + \rho g \\ \partial_t p_f + v \cdot \nabla p_f - \kappa \Delta p_f = 0 \end{array} \right.$$

where κ is the diffusion coefficient.

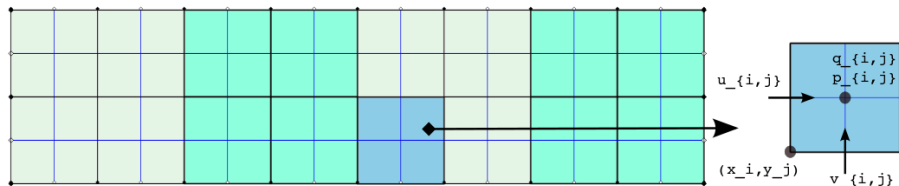
A photograph of a volcanic eruption. A large, dark, conical volcano is the central focus, with a massive, billowing plume of white ash and steam rising from its summit. The plume is dense and textured, filling much of the sky. The foreground shows a green, hilly landscape with some trees and a clear blue sky with light clouds. The overall scene is dramatic and powerful.

2 Numerical simulation

Dambreak: Mesh



Dambreak: how to treat the density?



$$\partial_t \rho + \mathbf{v} \cdot \nabla \rho = 0$$

numerical method:

RK3 TVD - WENO5 scheme.

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

We compute $(\tilde{v}^{n+1}, \Sigma^{n+1})$ solution of

$$\begin{cases} \frac{3\tilde{v}^{n+1} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n+1}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n+1}}{\rho^{n+1}} - \mathbf{e}_y, \\ \Sigma^{n+1} = \mathbb{P}_{q^n}(\Sigma^{n+1} + r D\tilde{v}^{n+1} + \varepsilon(\Sigma^n - \Sigma^{n+1})), \\ \tilde{v}^{n+1} \Big|_{\partial\Omega} = \mathbf{0}. \end{cases}$$

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

We compute $(\tilde{v}^{n+1}, \Sigma^{n+1})$ solution of

$$\begin{cases} \frac{3\tilde{v}^{n+1} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n+1}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n+1}}{\rho^{n+1}} - \mathbf{e}_y, \\ \Sigma^{n+1} = \mathbb{P}_{q^n}(\Sigma^{n+1} + r D\tilde{v}^{n+1} + \varepsilon(\Sigma^n - \Sigma^{n+1})), \\ \tilde{v}^{n+1} \Big|_{\partial\Omega} = \mathbf{0}. \end{cases}$$

We compute (v^{n+1}, p^{n+1}) thanks to the incompressibility constrain

$$\begin{cases} \frac{v^{n+1} - \tilde{v}^{n+1}}{\delta t} + \frac{2}{3\rho^{n+1}} \nabla(p^{n+1} - p^n) = \mathbf{0}, \\ \operatorname{div} v^{n+1} = 0, \quad v^{n+1} \cdot \mathbf{n} \Big|_{\partial\Omega} = 0. \end{cases}$$

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

We compute $(\tilde{v}^{n+1}, \Sigma^{n+1})$ solution of

$$\begin{cases} \frac{3\tilde{v}^{n+1} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n+1}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n+1}}{\rho^{n+1}} - \mathbf{e}_y, \\ \Sigma^{n+1} = \mathbb{P}_{q^n}(\Sigma^{n+1} + r D\tilde{v}^{n+1} + \varepsilon(\Sigma^n - \Sigma^{n+1})), \\ \tilde{v}^{n+1} \Big|_{\partial\Omega} = \mathbf{0}. \end{cases}$$

We compute (v^{n+1}, p^{n+1}) thanks to the incompressibility constrain

$$\begin{cases} \frac{v^{n+1} - \tilde{v}^{n+1}}{\delta t} + \frac{2}{3\rho^{n+1}} \nabla(p^{n+1} - p^n) = \mathbf{0}, \\ \operatorname{div} v^{n+1} = 0, \quad v^{n+1} \cdot \mathbf{n} \Big|_{\partial\Omega} = 0. \end{cases}$$

We compute p_f^{n+1} solution of

$$\frac{3p_f^{n+1} - 4p_f^n + p_f^{n-1}}{2\delta t} + 2v^{n+1} \cdot \nabla p_f^n - v^{n+1} \cdot \nabla p_f^{n-1} - \Delta p_f^{n+1} = 0.$$

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

We compute $(\tilde{v}^{n+1}, \Sigma^{n+1})$ solution of

$$\begin{cases} \frac{3\tilde{v}^{n+1} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n+1}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n+1}}{\rho^{n+1}} - \mathbf{e}_y, \\ \Sigma^{n+1} = \mathbb{P}_{q^n}(\Sigma^{n+1} + r D\tilde{v}^{n+1} + \varepsilon(\Sigma^n - \Sigma^{n+1})), \\ \tilde{v}^{n+1} \Big|_{\partial\Omega} = \mathbf{0}. \end{cases}$$

$\hookrightarrow k = 0$ $\Sigma^{n,0} = \Sigma^n$, and (v^n, p^n, p_f^n) given.

$\hookrightarrow k \geq 0$ $\Sigma^{n,k}$ known, we first compute $\tilde{v}^{n,k}$ solution of a

Laplace-type problem then we project the stress tensor to obtain $\Sigma^{n,k+1}$:

$$\begin{cases} \frac{3\tilde{v}^{n,k} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n,k}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n,k}}{\rho^{n+1}} - \mathbf{e}_y, \\ \tilde{v}^{n,k} \Big|_{\partial\Omega} = \mathbf{0}, \\ \Sigma^{n,k+1} = \mathbb{P}_{q^n}(\Sigma^{n,k} + r D\tilde{v}^{n,k} + \varepsilon(\Sigma^n - \Sigma^{n,k})). \end{cases}$$

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

We compute $(\tilde{v}^{n+1}, \Sigma^{n+1})$ solution of

$$\begin{cases} \frac{3\tilde{v}^{n+1} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n+1}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n+1}}{\rho^{n+1}} - \mathbf{e}_y, \\ \Sigma^{n+1} = \mathbb{P}_{q^n}(\Sigma^{n+1} + r D\tilde{v}^{n+1} + \varepsilon(\Sigma^n - \Sigma^{n+1})), \\ \tilde{v}^{n+1} \Big|_{\partial\Omega} = \mathbf{0}. \end{cases}$$

$\hookrightarrow k = 0$ $\Sigma^{n,0} = \Sigma^n$, and (v^n, p^n, p_f^n) given.

$\hookrightarrow k \geq 0$ $\Sigma^{n,k}$ known, we first compute $\tilde{v}^{n,k}$ solution of a

Laplace-type problem then we project the stress tensor to obtain $\Sigma^{n,k+1}$:

$$\begin{cases} \frac{3\tilde{v}^{n,k} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n,k}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n,k}}{\rho^{n+1}} - \mathbf{e}_y, \\ \tilde{v}^{n,k} \Big|_{\partial\Omega} = \mathbf{0}, \\ \Sigma^{n,k+1} = \mathbb{P}_{q^n}(\Sigma^{n,k} + r D\tilde{v}^{n,k} + \varepsilon(\Sigma^n - \Sigma^{n,k})). \end{cases}$$

Numerical scheme (without density)

$n = 0$ v^0, Σ^0, p^0 and p_f^0 given.

$n \geq 0$ v^n, Σ^n, p^n and p_f^n being calculated.

We compute $(\tilde{v}^{n+1}, \Sigma^{n+1})$ solution of

$$\begin{cases} \frac{3\tilde{v}^{n+1} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n+1}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n+1}}{\rho^{n+1}} - \mathbf{e}_y, \\ \Sigma^{n+1} = \mathbb{P}_{q^n}(\Sigma^{n+1} + r D\tilde{v}^{n+1} + \varepsilon(\Sigma^n - \Sigma^{n+1})), \\ \tilde{v}^{n+1} \Big|_{\partial\Omega} = \mathbf{0}. \end{cases}$$

$\hookrightarrow k = 0$ $\Sigma^{n,0} = \Sigma^n$, and (v^n, p^n, p_f^n) given.

$\hookrightarrow k \geq 0$ $\Sigma^{n,k}$ known, we first compute $\tilde{v}^{n,k}$ solution of a

Laplace-type problem then we project the stress tensor to obtain $\Sigma^{n,k+1}$:

$$\begin{cases} \frac{3\tilde{v}^{n,k} - 4v^n + v^{n-1}}{2\delta t} + 2v^n \cdot \nabla v^n - v^{n-1} \cdot \nabla v^{n-1} - \frac{\Delta \tilde{v}^{n,k}}{\rho^{n+1}} + \frac{\nabla p^n}{\rho^{n+1}} = \frac{\operatorname{div} \Sigma^{n,k}}{\rho^{n+1}} - \mathbf{e}_y, \\ \tilde{v}^{n,k} \Big|_{\partial\Omega} = \mathbf{0}, \\ \Sigma^{n,k+1} = \mathbb{P}_{q^n}(\Sigma^{n,k} + r D\tilde{v}^{n,k} + \varepsilon(\Sigma^n - \Sigma^{n,k})). \end{cases}$$

Experimental conditions

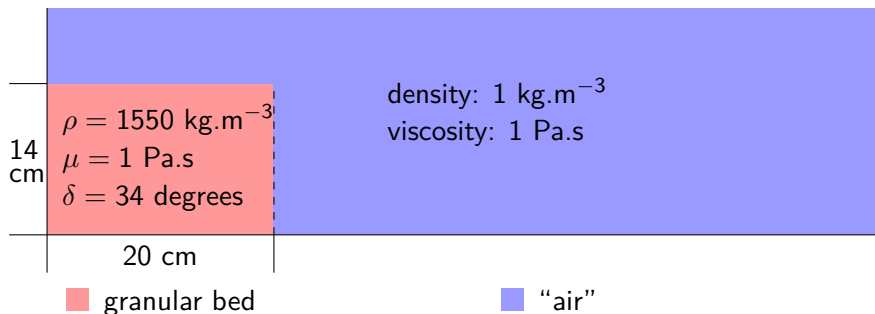


Figure: Experimental setup

Numerical simulation

We compute the collapse with different diffusion coefficient κ .

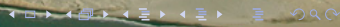
- With $\kappa = 1$, it happens nothing.
- $\kappa = 0.2$
- $\kappa = 0.1$

Diffusion coefficient = 0.2

Diffusion coefficient = 0.1



3 Perspectives



Perspectives

- To compare the results given by this numerical scheme for the **non-fluidised** case:
 - ▶ Lower yield stress, . . .
 - ▶ Lower friction coefficient. . . .
- To compare the results given by this numerical scheme for the **fluidised** case:
 - ▶ Adapt the viscosity value.

Thank you!

